## Fakeons and quantum gravity

Damiano Anselmi

· FLRW solution

· Cosmology

The problem of quantum gravity is to  
make it renormalizable and unitary at  
the same time  

$$\frac{1}{(p^2-m_1^2)(p^2-m_2^2)} = \frac{1}{m_1^2-m_2^2} \left[ \frac{1}{p^2-m_1^2} - \frac{1}{p^2-m_2^2} \right]$$
Higher derivatives lead to renormalizability,  
but violate unitarity, unless...

Unitarity: 
$$S^{\dagger}S = 1$$
  
 $S = A + iT$   
 $2 \operatorname{Im}T = T^{\dagger}T$  Optical theorem  
 $\frac{1}{2i} [\langle b|T|a \rangle - \langle b|T^{\dagger}|a \rangle] = \sum_{|n\rangle \in W} \langle b|T^{\dagger}|n \rangle \langle n|T|a \rangle, \quad |a \rangle, |b \rangle \in W$ 

In general, W may contain unphysical states  
in higher-derivative theories (those with 
$$\sigma_n = 1$$
)  
Feynman prescription  $\rightarrow$  pseudoumitarity equation

$$2\operatorname{Im}\left[\left(-i\right)\right) - \left\langle \right] = \left\langle \right\rangle + \left\langle \right| = \int d\Pi_{f} \left| \right\rangle - \left|^{2} \right\rangle$$
$$2\operatorname{Im}\left[\left(-i\right) - \left(-i\right)\right] = -\left(-i\right) - \left(-i\right) - \left(-i\right)\right|^{2} = -\left(-i\right) - \left(-i\right) - \left($$

Culting equations 
$$Propagator ; \qquad G(p,m) = \frac{1}{p^2 - m^2}$$

The Feynman prescription gives

$$G_{+}(p,m,\epsilon) = \frac{1}{p^2 - m^2 + i\epsilon}$$

The optical theorem is ok:

$$2\mathrm{Im}\left[(-i)\right] = \left| \left| \left| \left| \right|^{2} \right|^{2} \right|^{2} \right|^{2} \geq \mathcal{O}$$

Indeed :  $\operatorname{Im}\left[-\frac{1}{p^2 - m^2 + i\epsilon}\right] = \pi\delta(p^2 - m^2) \qquad \gtrless \bigcirc$ Gabost: opposite residue  $-\frac{1}{p^2 - m^2 + i\epsilon}$ The aptical theorem is violated Note that the prescription is crucial



This is adviewed by inserting on infinitesimal width as follows:

$$\mathbb{G}_{\pm}(p,m,\mathcal{E}^2) = \pm \frac{p^2 - m^2}{(p^2 - m^2)^2 + \mathcal{E}^4}$$

Note that the residue is zero on shell : => NO PARTICLE

Example : bubble diagram  

$$i\mathcal{M} \propto \int \frac{d\mathbf{k}^{o}}{2\pi} \int_{\mathbb{R}^{3}} \frac{d^{3}\mathbf{k}}{(2\pi)^{3}} \mathbb{G}_{\pm}(p-k,m_{1},\mathcal{E}^{2}) \mathbb{G}_{\pm}(k,m_{2},\mathcal{E}^{2})$$
  
 $\downarrow \mathcal{M}$   
 $= \int_{\mathbb{R}^{3}} \frac{d^{3}\mathbf{k}}{(2\pi)^{3}} w(p,\mathbf{k})$   
 $w(\mathbf{p},\mathbf{k})$  is singular for  
 $|p^{0}| = \sqrt{(\mathbf{p}-\mathbf{k})^{2} + m_{1}^{2} \pm i\mathcal{E}^{2}} + \sqrt{\mathbf{k}^{2} + m_{2}^{2} \pm i\mathcal{E}^{2}}$   
Here Lorentz invariance  $\mathcal{K}$  analiticity are violated





Since the prescription is symmetric w.r.t. the real axis, the imaginary part of the amplitude vanishes:  $0 = 2 \operatorname{Im} \left[ (-i) - \bigcirc^{\mathsf{F}} \right] = - \bigcirc^{\mathsf{F}} = \int^{\mathsf{F}} d\Pi_f \left| - \bigvee^{\mathsf{Im}} \right|^2 = \mathcal{O}$ => the fakeon F MUST be projected away

From  

$$\frac{1}{2i} \left[ \langle b | T | a \rangle - \langle b | T^{\dagger} | a \rangle \right] = \sum_{|n\rangle \in W} \langle b | T^{\dagger} | n \rangle \langle n | T | a \rangle, \quad |a \rangle, |b \rangle \in W$$
to  

$$\frac{1}{2i} \left[ \langle b | T | a \rangle - \langle b | T^{\dagger} | a \rangle \right] = \sum_{|n\rangle \in V} \langle b | T^{\dagger} | n \rangle \langle n | T | a \rangle, \quad |a \rangle, |b \rangle \in V$$
with  $V \subset W$ ,  $|F \rangle \in W$ ,  $|F \rangle \notin V$   

$$\left[ \ln \left( q^{2} - m^{2} + i \varepsilon \right) - \frac{1}{2} \left[ \ln \left( q^{2} - m^{2} \right)^{2} \right]$$
To all orders:  
D. Anselmi (2018), Fakeons & Lee-Wick models,  $]HEP$ 

D. Anseimi and M. Piva, Quantum gravity, fakeons and microcausality, J. High E. Phys. 11 (2018) 21, 18A3 Renormalization.com and arXiv:1806.03605 [hep-th]. igh Energy

$$\mathcal{S}_{\rm QG}(g,\phi,\chi,\Phi) = S_{\rm H}(g) + S_{\chi}(g,\chi) + S_{\phi}(\tilde{g},\phi) + S_{\mathfrak{m}}(\tilde{g}e^{\kappa\phi},\Phi)$$

Where 
$$\tilde{g}_{\mu\nu} = g_{\mu\nu} + 2\chi_{\mu\nu}$$
  
 $S_{\rm H}(g) = -\frac{\zeta}{2\kappa^2} \int \sqrt{-g}R, \qquad S_{\phi}(g,\phi) = \frac{3\zeta}{4} \int \sqrt{-g} \left[ \nabla_{\mu}\phi \nabla^{\mu}\phi - \frac{m_{\phi}^2}{\kappa^2} \left(1 - e^{\kappa\phi}\right)^2 \right]$   
 $S_{\chi}(g,\chi) = S_{\rm H}(\tilde{g}) - S_{\rm H}(g) - 2 \int \chi_{\mu\nu} \frac{\delta S_{\rm H}(\tilde{g})}{\delta g_{\mu\nu}} + \frac{\zeta^2}{2\alpha\kappa^2} \int \sqrt{-g} (\chi_{\mu\nu}\chi^{\mu\nu} - \chi^2) \Big|_{g \to \tilde{g}}.$ 

Now, the  $\chi_{\mu\nu}$  sction has the wrong overall soon :  $S_{\chi}(g,\chi) = -\frac{\zeta}{\kappa^2} S_{\rm PF}(g,\chi,m_{\chi}^2) - \frac{\zeta}{2\kappa^2} \int \sqrt{-g} R^{\mu\nu} (\chi \chi_{\mu\nu} - 2\chi_{\mu\rho} \chi_{\nu}^{\rho}) + S_{\chi}^{(>2)}(g,\chi)$ where  $S_{\rm PF}$  is the covariantized Pauli-Fiers action

This means that 
$$\chi_{\mu\nu}$$
 MUST be quantized as  
a fakeon. This way, we have both renormalizability  
and unitarity.



The fields of the standard model and quantum gravity could explain everything we know

Graviton multiplet: Shyn, &, Xnu &  $g_{\mu\nu} = M_{\mu\nu} + 2\kappa h_{\mu\nu}$ / spin-2 Jution of the metric m nassive fakeon of scalar mass mx massive Fakeon width:  $N_s + 6 N_f + 12 N_v$  $\Gamma_{\chi} = -\alpha_{\chi} C m_{\chi}$ Tx <O: causality is violated by Xnv  $\alpha_{\chi} = \left(\frac{M_{\chi}}{M_{\rm ex}}\right)^2$ 

D. A. and M. Piva, The ultraviolet behavior of quantum gravity, J. High Energy Phys. 05 (2018) 027 and arxiv:1803.0777 [hep-th]

JHEP 11 (2018) 21 D. A. and M. Piva, Quantum gravity, fakeons and microcausality, arxiv:1806.03605 [hep-th]

## PROJECTION

At the level of generating functionals:

 $\Gamma(\varphi,\chi)$ 

X = fakeons

Solve

 $\delta\Gamma(arphi,\chi)/\delta\chi\,=\,0$   $\,\,$  by means of the fakeon prescription

Let  $\langle X \rangle$  denote the solution

 $\varphi = physical fields$ 

Projected functionals:

$$\Gamma_{\rm pr}(\varphi) = \Gamma(\varphi, \langle \chi \rangle)$$

$$\begin{split} Z_{\rm pr}(J) &= \int [{\rm d}\varphi {\rm d}\chi] \exp\left(iS(\varphi,\chi) + i\int J\varphi\right) = \exp\left(iW_{\rm pr}(J)\right) \\ & \text{No source } \mathcal{J}_{\mathbf{x}} \text{ for } \mathbf{x} \end{split}$$

Projection = integrating out the fakeons with the fakeon prescription

Classical limit

## The action

$$S_{\rm QG}(g,\phi,\chi,\Phi) = S_{\rm H}(g) + S_{\chi}(g,\chi) + S_{\phi}(\tilde{g},\phi) + S_{\mathfrak{m}}(\tilde{g}e^{\kappa\phi},\Phi)$$
  
is NOT the classical limit, because it  
is unprojected  
Unprojected field equations:  
$$g_{\mu\nu} : \qquad R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R = \frac{\kappa^2}{\zeta} \left[e^{3\kappa\phi}fT^{\mu\nu}_{\mathfrak{m}}(\tilde{g}e^{\kappa\phi},\Phi) + fT^{\mu\nu}_{\phi}(\tilde{g},\phi) + T^{\mu\nu}_{\chi}(g,\chi)\right]$$

$$\begin{split} \varphi &: \quad -\frac{1}{\sqrt{-\tilde{g}}} \; \partial_{\mu} \left( \sqrt{-\tilde{g}} \tilde{g}^{\mu\nu} \partial_{\nu} \phi \right) - \frac{-\varphi}{\kappa} \left( \mathrm{e}^{\kappa\phi} - 1 \right) \mathrm{e}^{\kappa\phi} = \frac{1}{3\zeta} T^{\mu\nu}_{\mathfrak{m}} (\tilde{g} \mathrm{e}^{\kappa\phi}, \Phi) \tilde{g}_{\mu\nu}, \\ \chi_{\mu\nu} &: \qquad \frac{1}{\sqrt{-g}} \frac{\delta S_{\chi}(g, \chi)}{\delta \chi_{\mu\nu}} = \mathrm{e}^{3\kappa\phi} f T^{\mu\nu}_{\mathfrak{m}} (\tilde{g} \mathrm{e}^{\kappa\phi}, \Phi) + f T^{\mu\nu}_{\phi} (\tilde{g}, \phi), \end{split}$$

At the tree level, the subtleties about integration paths and average continuations are not important, so we can take  $\frac{p^2 - m^2}{(p^2 - m^2)^2 + \mathcal{E}^4} = \mathcal{P}\frac{1}{p^2 - m^2}$ 

Use it to solve the fakeon equations The projected classical action is then

 $\mathcal{S}_{\text{QG}} (g, \phi, \Phi) = S_{\text{H}}(g) + S_{\chi}(g, \langle \chi \rangle) + S_{\phi}(\bar{g}, \phi) + S_{\mathfrak{m}}(\bar{g}e^{\kappa\phi}, \Phi)$ 

where  $\langle \chi \rangle$  is the solution  $\bar{g}_{\mu\nu} = g_{\mu\nu} + 2 \langle \chi_{\mu\nu} \rangle$ 

Example : 
$$\mathcal{L}_{HD} = \frac{m}{2}(\dot{x}^2 - \tau^2 \ddot{x}^2) + xF_{ext}(t)$$
  
 $m \frac{d^2}{dt^2} \left(1 + \tau^2 \frac{d^2}{dt^2}\right) = F_{ext}$   
invert with  $\mathcal{P} \frac{1}{1 + \tau^2 \frac{d^2}{dt^2}}$   
The projected equation is  
 $m\ddot{x} = \int_{-\infty}^{\infty} du \frac{\sin(|u|/\tau)}{2\tau} F_{ext}(t-u)$   
 $\longrightarrow$  violation of microcousality  
Fma  $\langle F \rangle = ma$  !!

The FLRW metric

$$\Sigma = 1 + \frac{1}{m_{\phi}^2} \left( 3\frac{\dot{a}}{a} + \frac{\mathrm{d}}{\mathrm{d}t} \right) \frac{\mathrm{d}}{\mathrm{d}t}, \qquad \Upsilon = \Sigma + \frac{2}{m_{\phi}^2} \left[ \frac{k}{a^2} + 3\frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\dot{a}}{a} \right) \right].$$

$$\frac{\dot{a}^2}{a^2} + \frac{k}{a^2} = \frac{8\pi G}{3}\tilde{\rho},$$
$$2\frac{\ddot{a}}{a} + \frac{\dot{a}^2}{a^2} + \frac{k}{a^2} = -8\pi G\tilde{p}.$$

where

OR :

$$\begin{split} \tilde{\rho} &= \frac{1}{4} \langle \rho - 3p \rangle_{\Sigma} + \frac{3}{4} \langle \rho + p \rangle_{\Upsilon} \\ \tilde{p} &= \frac{1}{4} \langle \rho + p \rangle_{\Upsilon} - \frac{1}{4} \langle \rho - 3p \rangle_{\Sigma} \end{split}$$

The projection can be handled exactly for radiation combined with the vacuum energy:  

$$p = \frac{\rho}{3} + p_0 \qquad p_o = \text{constant}$$

Exact solution:  

$$\rho_0 = 3\sigma^2/(8\pi G) \quad \tilde{p} = (\tilde{\rho} - 4\rho_0)/3$$

$$\sigma_1 \sigma' = \text{constants}$$

$$\rho(t) = \frac{3}{8\pi G} \left(\sigma^2 + \frac{\sigma''^2}{4a^4}\right), \quad \sigma''^2 = \sigma'^2 \left(1 + \frac{4\sigma^2}{m_\phi^2}\right) \quad \tilde{\rho}(t) = \frac{3}{8\pi G} \left(\sigma^2 + \frac{\sigma'^2}{4a^4}\right)$$

$$a(t) = \sqrt{\frac{\sinh(\sigma t)}{\sigma}} \left(\sigma' \cosh(\sigma t) - \frac{k}{\sigma} \sinh(\sigma t)\right)$$

But in general the projection is defined parturbatively  
(since it comes from quantum gravity, which is defined  
parturbatively)  

$$\longrightarrow$$
 The classical equations are defined  
perturbatively : one may have to face  
asymptotic sories and nonperturbative effects  
(just to write the equations)  
Example: cosmic dust (p = 0) or p=MP  
 $NS = \frac{4}{3} - 1$ 

Is H 2 fokeon? Maybe...



D.A. On the nature of the Higgs boson, MPLA 33 (2019) 1950123 arXiv: 1811.02600 [hep-ph]



Comments on alternative approaches to the problem of quantum gravity

--- **string theory** is criticized for being nonpredictive. Moreover, its calculations often require mathematics that is not completely understood

--- **loop quantum gravity** is even more challenging, because it is at an earlier stage of development

--- holography (AdS/CFT correspondence) do not admit a weakly coupled expansion

--- asymptotic safety in nonperturbative as well.

None is as close to the standard model as the solution based on the fakeon idea, which is a quantum field theory, admits a perturbative expansion in terms of Feynman diagrams and allows us to make calculations with a comparable effort. It can be coupled to th standard model straightforwardly.

Our solution bests its competitors in calculatibity, predictivity and falsifiability. It is also rather rigid, because it contains only two new parameters. It could turn out to be the most predictive theory ever.