Fakeons, quantum gravity and the classical limit

- · How the idea of fake particle ("fakeon") leads to what I claim to be the right theory of quantum gravity
- · Physical implications and preductions
- · Classical limit
- · FLRW solution
- · Sosle invariance? Wont for it ...

Unitarity:

$$S^{\dagger}S = 1$$

$$S = 1 + iT$$

$$2\mathrm{Im}T = T^{\dagger}T$$

Optical theorem

In general, W may contain unphysical states

Diagrammatic version of the optical theorem:

$$2\operatorname{Im}\left[(-i)\right] = \left| \left| \left| \right| = \int d\Pi_f \right| \right|^2$$

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Culting equations

Propagator:
$$G(p,m) = \frac{1}{p^2 - m^2}$$
 A prescription is needed

The Feynman $G_{+}(p,m,\epsilon) = \frac{1}{p^2 - m^2 + i\epsilon}$ prescription gives

The optical theorem is ok:

$$2\operatorname{Im}\left[(-i)\right] = \left| \left\langle -i \right\rangle \right|^2 \geqslant 0$$

Indeed:

$$\operatorname{Im}\left[-\frac{1}{p^2 - m^2 + i\epsilon}\right] = \pi\delta(p^2 - m^2) \geqslant \bigcirc$$

The aptical theorem is violated Note that the prescription is crucial If we flip the sign of iE and the overall sopn, the optical theorem is ok:

$$\operatorname{Im}\left[\frac{1}{p^2 - m^2 - i\epsilon}\right] = \pi\delta(p^2 - m^2) \geqslant 0$$

Problem:

$$G_{\pm}(p, m, \epsilon) = \pm \frac{1}{p^2 - m^2 \pm i\epsilon}$$

cannot coexist

S.

U.G. Aglietti and D. Anselmi, Inconsistency of Minkowski higher-derivative theories, Eur. Phys. J. C 77 (2017) 84, 16A2 Renormalization.com and arXiv:1612.06510 [hep-th] The fakeon is the answer $G(p,m) = \frac{1}{p^2 - m^2}$

Write $\pm \frac{p^2 - m^2}{(p^2 - m^2)^2}$ and split the poles into pairs

$$\begin{array}{c|c} & & & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & \\ & & & \\ & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ &$$

This is advised by inserting on infinitesimal width as follows:

$$\mathbb{G}_{\pm}(p, m, \mathcal{E}^2) = \pm \frac{p^2 - m^2}{(p^2 - m^2)^2 + \mathcal{E}^4} = \pm \frac{1}{2} \left[G_+(p, m, \mathcal{E}^2) - G_-(p, m, \mathcal{E}^2) \right]$$

Note that the residue is zero on shell:

=> NO PARTICLE

BUT: what about the bad divergences?
Those occur in Minkowski space
which is not what we are doing here

Example: bubble diagram

$$\mathcal{U} \propto \int \frac{d\mathbf{k}}{2\pi} \int_{\mathbb{R}^3} \frac{d^3\mathbf{k}}{(2\pi)^3} \mathbb{G}_{\pm}(p-k, m_1, \mathcal{E}^2) \mathbb{G}_{\pm}(k, m_2, \mathcal{E}^2)$$

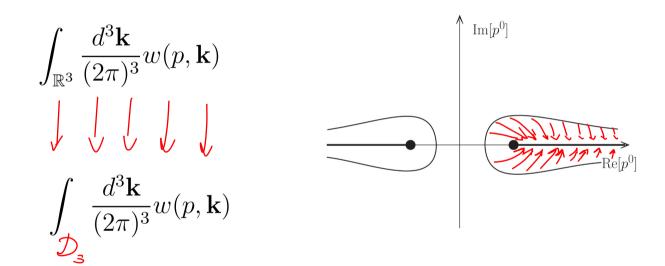
$$= \int_{\mathbb{R}^3} \frac{d^3\mathbf{k}}{(2\pi)^3} w(p, \mathbf{k})$$

 $|p^{0}| = \sqrt{(\mathbf{p} - \mathbf{k})^{2} + m_{1}^{2} \pm i\mathcal{E}^{2}} + \sqrt{\mathbf{k}^{2} + m_{2}^{2} \pm i\mathcal{E}^{2}}$

Here Lorentz invariance Ranaliticity are violated

 $\operatorname{Im}[p^0]$

The integration domain on the loop space momentum must be deformed



Lorentz invariance and analityaity are recovered in the limit where the areas are shrinked to branch outs

Re
$$[p^0]$$
 $\operatorname{Im}[p^0]$

Euclidean region

Region

above the fakeon threshold

 $\operatorname{Re}[p^0]$

Since the prescription is symmetric w.r.t. the real exis, the imaginary part of the amplitude vanishes: $0 = 2 \text{Im} \left[(-i) - \bigcap_{i=1}^{F} -$

$$0 = 2 \text{Im} \left[(-i) \longrightarrow_{\mathsf{F}} \right] = \longrightarrow_{\mathsf{F}} = \iint_{\mathsf{F}} \mathrm{d}\Pi_f \left| - \left\langle \right| = \right.$$

$$=) \text{ the falseon } \mathsf{F} \text{ MUST be projected away} ;$$

From

$$\frac{1}{2i} \left[\langle b|T|a \rangle - \langle b|T^{\dagger}|a \rangle \right] = \sum_{|n\rangle \in W} \langle b|T^{\dagger}|n\rangle \langle n|T|a \rangle, \qquad |a\rangle, |b\rangle \in W$$

to

$$\frac{1}{2i} \left[\langle b|T|a \rangle - \langle b|T^{\dagger}|a \rangle \right] = \sum \langle b|T^{\dagger}|n \rangle \langle n|T|a \rangle, \qquad |a \rangle, |b \rangle \in V$$

with
$$V \subset W$$
, $|F\rangle \in W$, $|F\rangle \notin V$

$$|n(q^2 - m^2 + i\varepsilon) \rightarrow |n(q^2 - m^2)^2|$$

To all orders:

D. Anselmi (2018), Fakeons & Lee-Wick models, JHEP

Quantum gravity

D. Anselmi, On the quantum field theory of the gravitational interactions, J. High Energy Phys. 06 (2017) 086, 17A3 Renormalization.com and arXiv: 1704.07728 [hep-th].

Consider the (renormalizable) higher-derivative action

$$S_{\rm QG} = -\frac{1}{2\kappa^2} \int \sqrt{-g} \left[2\Lambda_C + \zeta R + \alpha \left(R_{\mu\nu} R^{\mu\nu} - \frac{1}{3} R^2 \right) - \frac{\xi}{6} R^2 \right] + S_{\mathfrak{m}}(g, \Phi)$$

Eliminate the higher derivatives by means of

extra fields;

D. Anselmi and M. Piva, Quantum gravity, fakeons and microcausality, J. High Energy Phys. 11 (2018) 21, 18A3 Renormalization.com and arXiv:1806.03605 [hep-th].

$$S_{\rm QG}(g,\phi,\chi,\Phi) = S_{\rm H}(g) + S_{\chi}(g,\chi) + S_{\phi}(\tilde{g},\phi) + S_{\mathfrak{m}}(\tilde{g}e^{\kappa\phi},\Phi)$$

where $\tilde{g}_{\mu\nu}=g_{\mu\nu}+2\chi_{\mu\nu}$

$$S_{\rm H}(g) = -\frac{\zeta}{2\kappa^2} \int \sqrt{-g} R, \qquad S_{\phi}(g,\phi) = \frac{3\zeta}{4} \int \sqrt{-g} \left[\nabla_{\mu} \phi \nabla^{\mu} \phi - \frac{m_{\phi}^2}{\kappa^2} \left(1 - e^{\kappa \phi} \right)^2 \right]$$

 $S_{\chi}(g,\chi) = S_{\mathrm{H}}(\tilde{g}) - S_{\mathrm{H}}(g) - 2 \int \chi_{\mu\nu} \frac{\delta S_{\mathrm{H}}(\tilde{g})}{\delta g_{\mu\nu}} + \frac{\zeta^2}{2\alpha\kappa^2} \int \sqrt{-g} (\chi_{\mu\nu} \chi^{\mu\nu} - \chi^2) \big|_{g \to \tilde{g}}.$

Now, the $\chi_{\mu\nu}$ action has the wrong overall soon: $S_{\chi}(g,\chi) = -\frac{\zeta}{\kappa^2} S_{\rm PF}(g,\chi,m_{\chi}^2) - \frac{\zeta}{2\kappa^2} \int \sqrt{-g} R^{\mu\nu} (\chi \chi_{\mu\nu} - 2\chi_{\mu\rho} \chi_{\nu}^{\rho}) + S_{\chi}^{(>2)}(g,\chi)$ where $S_{\rm PF}$ is the covariantized Pauli-Fierz action

This means that $\chi_{\mu\nu}$ MUST be quantized as a fakeon. This way, we have both renormalizability and unitarity.

Instead, & can be quantized either as a fakeon or as a physical particle

Classical limit

The action

$$S_{\mathrm{QG}}(g,\phi,\chi,\Phi) = S_{\mathrm{H}}(g) + S_{\chi}(g,\chi) + S_{\phi}(\tilde{g},\phi) + S_{\mathfrak{m}}(\tilde{g}e^{\kappa\phi},\Phi)$$

is NOT the classical limit, because it

is unprojected

Unprojected field equations:

$$\mathcal{G}_{\mu\nu} : \qquad R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R = \frac{\kappa^2}{\zeta} \left[\mathrm{e}^{3\kappa\phi} f T^{\mu\nu}_{\mathfrak{m}} (\tilde{g} \mathrm{e}^{\kappa\phi}, \Phi) + f T^{\mu\nu}_{\phi} (\tilde{g}, \phi) + T^{\mu\nu}_{\chi} (g, \chi) \right]$$

$$\psi : \qquad -\frac{1}{\sqrt{-\tilde{g}}} \; \partial_{\mu} \left(\sqrt{-\tilde{g}} \tilde{g}^{\mu\nu} \partial_{\nu} \phi \right) - \frac{m_{\phi}^2}{\kappa} \left(\mathrm{e}^{\kappa\phi} - 1 \right) \mathrm{e}^{\kappa\phi} = \frac{\kappa \mathrm{e}^{3\kappa\phi}}{3\zeta} T^{\mu\nu}_{\mathfrak{m}} (\tilde{g} \mathrm{e}^{\kappa\phi}, \Phi) \tilde{g}_{\mu\nu},$$

$$\chi_{\mu\nu} : \qquad \frac{1}{\sqrt{-g}} \frac{\delta S_{\chi}(g, \chi)}{\delta \chi_{\mu\nu}} = \mathrm{e}^{3\kappa\phi} f T^{\mu\nu}_{\mathfrak{m}} (\tilde{g} \mathrm{e}^{\kappa\phi}, \Phi) + f T^{\mu\nu}_{\phi} (\tilde{g}, \phi),$$

At the tree level, the subtleties about integration paths and integration domains are not important,

So We can take
$$\frac{p^2 - m^2}{(p^2 - m^2)^2 + \mathcal{E}^4} = \mathcal{P} \frac{1}{p^2 - m^2}$$

Use it to solve the fakeon equations

The projected classical action is then

$$S_{\mathrm{QG}}(g,\phi,\Phi) = S_{\mathrm{H}}(g) + S_{\chi}(g,\langle\chi\rangle) + S_{\phi}(\bar{g},\phi) + S_{\mathfrak{m}}(\bar{g}e^{\kappa\phi},\Phi)$$

$$\bar{g}_{\mu\nu} = g_{\mu\nu} + 2\langle \chi_{\mu\nu} \rangle$$

Example:

$$\mathcal{L}_{HD} = \frac{m}{2}(\dot{x}^2 - \tau^2 \ddot{x}^2) + x F_{\text{ext}}(t)$$

$$m \frac{d^2}{dt^2} \left(1 + \tau^2 \frac{d^2}{dt^2}\right) \times = \overline{T}_{ext}$$

The projected equation is

$$m\ddot{x} = \int_{-\infty}^{\infty} du \frac{\sin(|u|/\tau)}{2\tau} F_{\text{ext}}(t-u)$$



Unprojected equations
$$\left(\frac{\mathcal{L}}{\sqrt{-g}} \sim R + R^2\right)$$

$$\Sigma\left(\frac{\ddot{a}}{a} + \frac{\dot{a}^2}{a^2} + \frac{k}{a^2}\right) = \frac{4\pi G}{3}(\rho - 3p), \qquad \Upsilon\left(\frac{\ddot{a}}{a} - \frac{\dot{a}^2}{a^2} - \frac{k}{a^2}\right) = -4\pi G(\rho + p),$$

$$\Sigma(1 - e^{-\kappa\phi}) = -\frac{8\pi G}{3m_{\phi}^2}(\rho - 3p)$$

where

$$\Sigma = 1 + \frac{1}{m_{\phi}^2} \left(3\frac{\dot{a}}{a} + \frac{\mathrm{d}}{\mathrm{d}t} \right) \frac{\mathrm{d}}{\mathrm{d}t}, \qquad \Upsilon = \Sigma + \frac{2}{m_{\phi}^2} \left[\frac{k}{a^2} + 3\frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{\dot{a}}{a} \right) \right].$$

Projection. Define $\langle A \rangle_X \equiv \frac{1}{2} \left| \frac{1}{X} \right|_{\text{rit}} + \left| \frac{1}{X} \right|_{\text{adv}} A$ and use it to define 1 and 1

$$1-\mathrm{e}^{-\kappa\phi}=-\frac{8\pi G}{3m_\phi^2}\langle\rho-3p\rangle_\Sigma$$
 The projection can be handled exactly for radiation combained with the vacuum energy:

 $p = \frac{\rho}{3} + p_0$

 $\frac{\ddot{a}}{a} + \frac{\dot{a}^2}{a^2} + \frac{k}{a^2} = \frac{4\pi G}{3} \langle \rho - 3p \rangle_{\Sigma}$

 $\frac{\ddot{a}}{a} - \frac{\dot{a}^2}{a^2} - \frac{k}{a^2} = -4\pi G \langle \rho + p \rangle_{\Upsilon}$

it is NOT over yet

P. = constant

Conclusions

Quantum field theory of particles and fakeons offers new opportunities to high-energy physics and quantum gravity

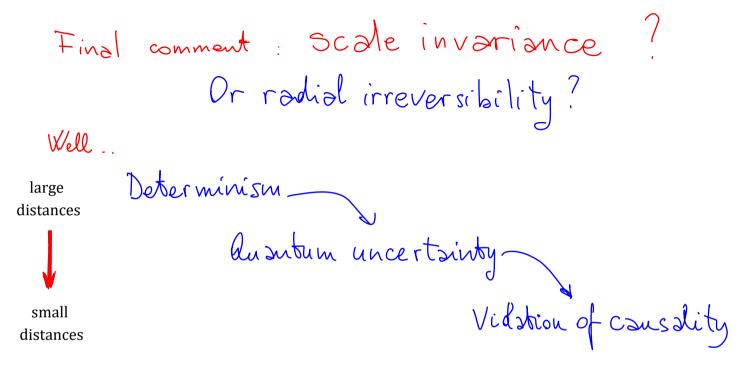
Fakeons cannot be observed directly. Their presence can however be detected by means of precision measurements, starting from those that probe the contributions coming from the imaginary parts of loop diagrams, above the fakeon thresholds

Higgs?

Fakeons allow us to quantize gravity and lead to the violation of causality at energies larger than their masses

The classical action is just an interim, local action that must be projected to obtain the true classical limit.

The classicization is perturbative: asymptotic series



It appears that there is a pattern:

Chains up here, freedom (anarchy?) down there

The universe is ... un-inside-out-able

