

One-loop Effective Action Of Quantum Gravity With Fakeons And Its Phenomenology

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Modena, September 9th, 2019

Avenues of Quantum Field Theory in Curved Spacetime

The problem of renormalizability in QG

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Hilbert–Einstein action: unitary but nonrenormalizable theory

$$S_{\text{HE}} = -\frac{1}{2\kappa^2} \int \sqrt{-g} R, \quad \kappa^2 = 8\pi G.$$

$$\Gamma_{\text{HE}}^{(\text{ct})} = -\frac{1}{2\kappa^2} \int \sqrt{-g} [c_1 R^2 + c_2 R_{\mu\nu} R^{\mu\nu} + c_3 R_{\rho\sigma}^{\mu\nu} R_{\alpha\beta}^{\rho\sigma} R_{\mu\nu}^{\alpha\beta} + \underbrace{\dots}_{\infty}].$$

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Stelle action: renormalizable but not unitary theory

$$S_{\text{HD}} = -\frac{1}{2\kappa^2} \int \sqrt{-g} [\gamma R + \alpha R_{\mu\nu} R^{\mu\nu} + \beta R^2 + 2\Lambda_{\text{C}}].$$

$$\Gamma_{\text{HD}}^{(\text{ct})} = -\frac{1}{2\kappa^2} \int \sqrt{-g} [a_\gamma R + a_\alpha R_{\mu\nu} R^{\mu\nu} + a_\beta R^2 + 2a_{\text{C}}].$$

Unitarity

$$SS^\dagger = 1, \quad S = 1 + iT.$$

Unitarity equation (optical theorem)

$$2\text{Im}T = TT^\dagger.$$

Cutting equations

↓

Pseudo-unitarity equation

$$2\text{Im}T = THT^\dagger.$$

$$H = \text{diag}(\dots, 1, \dots, 1, -1, \dots, -1, \dots).$$

Prescription for the propagator $\frac{\pm 1}{k^2 - m^2}$

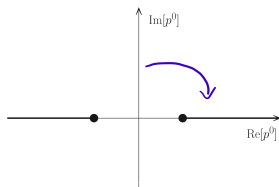
Prescription for the amplitude $\mathcal{A}(p)$

Prescription for the propagator $\frac{\pm 1}{k^2 - m^2}$

Standard particle : $\frac{1}{k^2 - m^2 + i\epsilon}$.

Ghost : $\frac{-1}{k^2 - m^2 + i\epsilon}$.

Prescription for the amplitude $\mathcal{A}(p)$



$$\mathcal{A}_+(p) = \mathcal{A}(p + i\epsilon)$$

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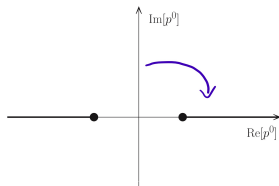
Fake particle (fakeon):

$$\pm \frac{k^2 - m^2}{(k^2 - m^2)^2 + \mathcal{E}^4}.$$

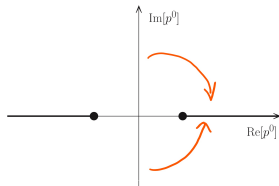
+ integration domain deformations

(see D. Anselmi and MP, JHEP 1706 (2017) 066.)

Prescription for the amplitude $\mathcal{A}(p)$



$$\mathcal{A}_+(p) = \mathcal{A}(p + i\epsilon)$$



$$\mathcal{A}_{\text{AV}}(p) = \frac{1}{2} [\mathcal{A}_+(p) + \mathcal{A}_-(p)].$$

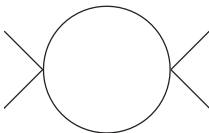
A theory of particles and fakeons is unitary.

D. Anselmi and MP, PRD 96 (2017) 045009.

D. Anselmi, JHEP 02 (2018) 141.

$$\mathcal{L} = \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - \frac{\lambda}{4!} \varphi^4.$$

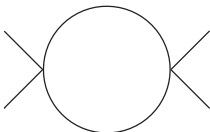
One-loop bubble diagram



Compute the Euclidean amplitude $\mathcal{A}(p)$ and then use the prescription.

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One-loop bubble diagram



Compute the Euclidean amplitude $\mathcal{A}(p)$ and then use the prescription.

After renormalizing the UV divergence

Feynman prescription

$$\mathcal{A}_+(p) = \frac{1}{2(4\pi)^2} \ln \frac{-p^2 - i\epsilon}{\mu^2}.$$

New prescription

$$\mathcal{A}_{AV}(p) = \frac{1}{4(4\pi)^2} \ln \frac{(p^2)^2}{\mu^4}.$$

The theory of quantum gravity and fakeons

D. Anselmi, JHEP 1706 (2017) 086.

$$S_{\text{HD}} = -\frac{1}{2\kappa^2} \int \sqrt{-g} \left[2\Lambda_C + \zeta R + \alpha \left(R_{\mu\nu} R^{\mu\nu} - \frac{1}{3} R^2 \right) - \frac{\xi}{6} R^2 \right].$$

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The propagator in De Donder gauge ($\Lambda_C = 0$, $\alpha = \xi$, $g_{\mu\nu} = \eta_{\mu\nu} + 2\kappa h_{\mu\nu}$) is

$$\langle h_{\mu\nu}(p) h_{\rho\sigma}(-p) \rangle_0 = \frac{i}{2p^2(\zeta - \alpha p^2)} \mathcal{I}_{\mu\nu\rho\sigma} = \left\{ \frac{1}{p^2} - \frac{1}{p^2 - \zeta/\alpha} \right\} \frac{i}{2\zeta} \mathcal{I}_{\mu\nu\rho\sigma},$$

$$\mathcal{I}_{\mu\nu\rho\sigma} \equiv (\eta_{\mu\rho}\eta_{\nu\sigma} + \eta_{\mu\sigma}\eta_{\nu\rho} - \eta_{\mu\nu}\eta_{\rho\sigma}).$$

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With the new prescription it turns into

$$\left\{ \frac{1}{p^2 + i\epsilon} - \frac{p^2 - \zeta/\alpha}{[(p^2 - \zeta/\alpha + i\epsilon)^2 + \mathcal{E}^4]_{\text{AV}}} \right\} \frac{i}{2\zeta} \mathcal{I}_{\mu\nu\rho\sigma}.$$

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Renormalization

The fakeon prescription does not spoil renormalizability.

Using the Batalin-Vilkovisky formalism we computed the beta functions in a new method by renormalizing the sources of the BRST tf's.

D. Anselmi and MP, JHEP 05 (2018) 027.

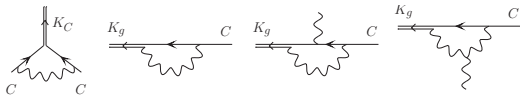
$$S_{\text{count}} = \frac{\mu^{-\varepsilon}}{(4\pi)^2 \varepsilon} \int \sqrt{-g} \left[2\Delta\Lambda_C + \Delta\zeta R + \Delta\alpha \left(R_{\mu\nu} R^{\mu\nu} - \frac{1}{3} R^2 \right) - \frac{\Delta\xi}{6} R^2 \right] + (S, \mathcal{F}),$$

$$\Delta\alpha = -\frac{133}{10}, \quad \Delta\xi = \frac{5}{6} + \frac{5\xi}{\alpha} + \frac{5\xi^2}{3\alpha^2}, \quad \Delta\zeta = \zeta \left(\frac{5}{6\xi} + \frac{5\xi}{3\alpha^2} + A \right),$$

$$\Delta\Lambda_C = \Lambda_C \left(-\frac{5}{\alpha} + \frac{2}{\xi} - 2A \right) - \frac{5\zeta^2}{4\alpha^2} - \frac{\zeta^2}{4\xi^2}.$$

The beta functions are

$$\beta_\alpha = -\frac{2\kappa^2}{(4\pi)^2} \Delta\alpha, \quad \beta_\xi = -\frac{2\kappa^2}{(4\pi)^2} \Delta\xi, \quad \beta_\zeta = -\frac{2\kappa^2}{(4\pi)^2} \Delta\zeta, \quad \beta_{\Lambda_C} = -\frac{2\kappa^2}{(4\pi)^2} \Delta\Lambda_C.$$



Absorptive part of graviton self energy (part I)

$\Lambda_C = 0$, massless matter

$$\mathcal{M}_{\text{abs}} = \text{Im}\mathcal{M} \quad \mathcal{D}_{\text{abs}} = -\text{Re}\mathcal{D}, \quad \mathcal{M} = \text{amplitude}, \quad \mathcal{D} = \text{diagram}.$$

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$$\langle h_{\mu\nu}(p)h_{\rho\sigma}(-p) \rangle_0 = \langle h_{\mu\nu}(p)h_{\rho\sigma}(-p) \rangle_{0\text{grav}} + \langle h_{\mu\nu}(p)h_{\rho\sigma}(-p) \rangle_{0\text{fake}}.$$

Possible cases in the diagram (pure gravity case)

- i) **Fake-Fake:** AV continuation \longrightarrow purely imaginary;

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Possible cases in the diagram (pure gravity case)

- i) **Fake-Fake:** AV continuation \rightarrow purely imaginary;
- ii) **Grav-Fake:** AV continuation \rightarrow purely imaginary;
- iii) **Grav-Grav:** Standard Wick rotation \rightarrow nontrivial real part.

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- i) **Fake-Fake:** AV continuation \rightarrow purely imaginary;
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The contributions of type iii) can be evaluated by using

$$\langle h_{\mu\nu}(p)h_{\rho\sigma}(-p) \rangle_0 = \langle h_{\mu\nu}(p)h_{\rho\sigma}(-p) \rangle_{0\text{grav}}.$$

In the case of pure gravity the absorptive part can be written as

$$\Gamma_{\text{abs}} = - \int \frac{\delta S_{\text{HD}}}{\delta g_{\mu\nu}} \Delta g_{\mu\nu}, \quad \Delta g_{\mu\nu} \text{ piecewise local function of } h_{\mu\nu}.$$

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Adding (massless) matter fields of all types the absorptive part is

$$\Gamma_{\text{abs}} = \frac{i\mu^{-\varepsilon}}{16\pi} \int \sqrt{-g} \left[C \left(R_{\mu\nu} \theta(-\square_c) R^{\mu\nu} - \frac{1}{3} R \theta(-\square_c) R \right) + \frac{N_s \eta^2}{36} R \theta(-\square_c) R \right] - \int \frac{\delta S_{\text{HD}}}{\delta g_{\mu\nu}} \Delta g_{\mu\nu},$$

$$C = \frac{1}{120} (N_s + 6N_f + 12N_v),$$

$\eta =$ nonminimal coupling for scalar fields.

Absorptive part of graviton self energy (part II)

$\Lambda_C = 0$, massive matter

D. Anselmi and MP, JHEP 11 (2018) 021.

Equivalent action:

auxiliary fields ϕ , $\chi_{\mu\nu}$ + Weyl transformation + field redefinitions.

Absorptive part of graviton self energy (part II)

$\Lambda_C = 0$, massive matter

D. Anselmi and MP, JHEP 11 (2018) 021.

Equivalent action:
auxiliary fields ϕ , $\chi_{\mu\nu}$ + Weyl transformation + field redefinitions.

$$S_{\text{QG}}(g, \phi, \chi, \Phi) = S_{\text{H}}(g) + S_{\chi}(g, \chi) + S_{\phi}(\tilde{g}, \phi) + S_m(\tilde{g}e^{\kappa\phi}, \Phi),$$
$$\tilde{g}_{\mu\nu} = g_{\mu\nu} + 2\chi_{\mu\nu}.$$

$$S_{\text{H}} = -\frac{\zeta}{2\kappa^2} \int \sqrt{-g} R, \quad S_{\phi}(g, \phi) = \frac{3\zeta}{4} \int \sqrt{-g} \left[\nabla_{\mu}\phi\nabla^{\mu}\phi - \frac{m_{\phi}^2}{\kappa^2} (1 - e^{\kappa\phi})^2 \right].$$

$$S_{\chi}(g, \chi) = S_{\text{H}}(\tilde{g}) - S_{\text{H}}(g) - 2 \int \chi_{\mu\nu} \frac{\delta S_{\text{H}}(\tilde{g})}{\delta g_{\mu\nu}} + \frac{\zeta^2}{2\alpha\kappa^2} \int \sqrt{-g} (\chi_{\mu\nu}\chi^{\mu\nu} - \chi^2)|_{g \rightarrow \tilde{g}}.$$

S_m = Standard Model,

Graviton multiplet: $G_A = \{h_{\mu\nu}, \phi, \chi_{\rho\sigma}\}$.

$\{G, F, F\}, \quad \{G, S, F\}, \quad \{G, S, Gh\}$.

Computation of the absorptive parts $M_{AB} = \langle G_A G_B \rangle_{\text{abs}}^{1\text{-loop}}$

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$$\Gamma_{\text{abs}}^{hh} = - \int \frac{\delta S_{\text{QG}}}{\delta g_{\mu\nu}} \Delta g_{\mu\nu},$$

$$\Gamma_{\text{abs}}^{\Phi} = \frac{i}{16\pi} \int \sqrt{-g} R_{\mu\nu} \theta(r_{\Phi}) \theta(1-r_{\Phi}) \sqrt{1-r_{\Phi}} \left[P_{\Phi}(r_{\Phi}) \left(R^{\mu\nu} - \frac{1}{3} g^{\mu\nu} R \right) + Q_{\Phi}(r_{\Phi}) g^{\mu\nu} R \right],$$

$$r_{\Phi} = -4m_{\Phi}^2/\square, \quad P_{\Phi}, Q_{\Phi} = \text{polynomials.}$$

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Contributions to the absorptive part

$$\Gamma_{\text{abs}}^{\text{GFF}} = \Gamma_{\text{abs}}^{hh} + \Gamma_{\text{abs}}^m,$$

$$\Gamma_{\text{abs}}^{\text{GSF}} = \Gamma_{\text{abs}}^{hh} + \Gamma_{\text{abs}}^{\phi h} + \Gamma_{\text{abs}}^{\phi\phi} + \Gamma_{\text{abs}}^m,$$

$$\Gamma_{\text{abs}}^{\text{GSGh}} = \Gamma_{\text{abs}}^{\text{GSF}} + \Gamma_{\text{abs}}^{\chi h} + \Gamma_{\text{abs}}^{\chi\phi} + \Gamma_{\text{abs}}^{\chi\chi}.$$

Violation of microcausality

Resumming the self energies the corrected χ propagator at the peak m_χ is

$$\langle \chi_{\mu\nu}(p)\chi_{\rho\sigma}(-p) \rangle_{s \sim \bar{m}_\chi^2} = -\frac{i\kappa^2}{\zeta} \frac{Z_\chi}{s - \bar{m}_\chi^2 + i\bar{m}_\chi\Gamma_\chi} \Pi_{\mu\nu\rho\sigma}^{(2)}(p, s),$$

$$\Gamma_\chi = -\frac{m_\chi^3}{M_{\text{Pl}}^2} C, \quad C = \frac{N_s + 6N_f + 12N_v}{120}, \quad \Gamma_\chi < 0.$$

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Breit-Wigner distribution

$$\frac{i}{E - m + i\frac{\Gamma}{2}} \longrightarrow \text{sgn}(t)\theta(\Gamma t)\exp\left(-imt - \frac{\Gamma t}{2}\right).$$

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$$\text{Duration} \sim 1/|\Gamma_\chi|$$

$$\text{If } m_\chi \sim 10^{11}\text{GeV then } 1/|\Gamma_\chi| \sim 7 \cdot 10^{-17}\text{s}$$

$$\text{If } m_\chi \sim 10^{12}\text{GeV then } 1/|\Gamma_\chi| \sim 4 \cdot 10^{-20}\text{s}$$

For time intervals of the order $1/|\Gamma_\chi|$ past, present and future, as well as cause and effect lose meaning.

$$2\text{Im} \left[i \langle \chi_{\mu\nu}(p) \chi_{\rho\sigma}(-p) \rangle_{s \sim \bar{m}_\chi^2} \right] \xrightarrow{\Gamma_\chi \rightarrow 0^-} \frac{2\pi\kappa^2}{\zeta} Z_\chi \delta(s - \bar{m}_\chi^2) \Pi_{\mu\nu\rho\sigma}^{(2)}(p, s) \geq 0.$$

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In the case of the Stelle theory (GSGh) at tree level

$$\begin{aligned} 2\text{Im} \left[i \langle \chi_{\mu\nu}(p) \chi_{\rho\sigma}(-p) \rangle_0 \right] &= \frac{2\kappa^2}{\zeta} \text{Im} \left[\frac{1}{p^2 - m_\chi^2 + i\epsilon} \right] \Pi_{\mu\nu\rho\sigma}^{(2)}(p, s) \\ &= -\frac{2\pi\kappa^2}{\zeta} \delta(s - m_\chi^2) \Pi_{\mu\nu\rho\sigma}^{(2)}(p, s) \leq 0. \end{aligned}$$

$$2\text{Im} \left[\text{Diagram} \right] \neq \int d\Pi \left| \text{Diagram} \right|^2$$

The classical limit

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Example: HD particle in an external force

$$m\ddot{x} + m\tau^2 \dddot{x} = \left(1 + \tau^2 \frac{d^2}{dt^2}\right) m\ddot{x} = F_{\text{ext}}(t).$$

↓

$$m\ddot{x} = \int_{-\infty}^{\infty} G_F(t-t') F_{\text{ext}}(t') dt' \equiv \langle F_{\text{ext}} \rangle.$$

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↓

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Analogously the classical limit of QG leads to

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 8\pi G \langle T_{\mu\nu} \rangle.$$

For the explicit case of FLRW metric see D. Anselmi, JHEP 1904 (2019) 061.

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Perturbative unitarity: the projection $W \rightarrow V$ is consistent

HD Harmonic oscillator

$$S(x, \dot{x}, \ddot{x}) = \int dt \left[\frac{1}{2} m (\dot{x}^2 - \tau^2 \ddot{x}^2 - \omega^2 x^2) \right].$$

↓

$$S'(q, \dot{q}, Q, \dot{Q}) = \frac{m}{2} \int dt \left[\dot{q}^2 - \dot{Q}^2 + \frac{Q^2}{\tau^2} - \omega^2 (q + Q)^2 \right].$$

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$$H'(p, q, P, Q) \rightarrow H_W(u, p_u, v, p_v) = \Omega \left(a_u^\dagger a_u - \frac{1}{2} \right) - \tilde{\Omega} \left(a_v^\dagger a_v + \frac{1}{2} \right),$$

$$\Omega = \frac{1}{\tau\sqrt{2}} \sqrt{1 - \sqrt{1 - 4\tau^2\omega^2}}, \quad \tilde{\Omega} = \frac{1}{\tau\sqrt{2}} \sqrt{1 + \sqrt{1 - 4\tau^2\omega^2}}.$$

Projected Hamiltonian

$$H_W(u, p_u, v, p_v) \longrightarrow H_V(u, p_u) = \Omega a_u^\dagger a_u + \frac{1}{2} (\Omega - \tilde{\Omega}).$$

Space fourier transform of quantum cosmological perturbations have same quadratic actions but with the substitutions

$$m \rightarrow m(t), \quad \omega \rightarrow \omega(t), \quad dt \rightarrow a(t)^3 dt.$$

Summary

- A new quantization prescription gives a unitary and renorm. QFT of gravity.
 - Predictive and possibly testable in the near future.
 - Quantization of cosmological perturbations (E. Bianchi and MP, in preparation).
- New type of degrees of freedom (fakeons).
- New phenomenology due to the presence of fakeons.